

ECE 107: Electromagnetism

Set 4: Maxwell's equations

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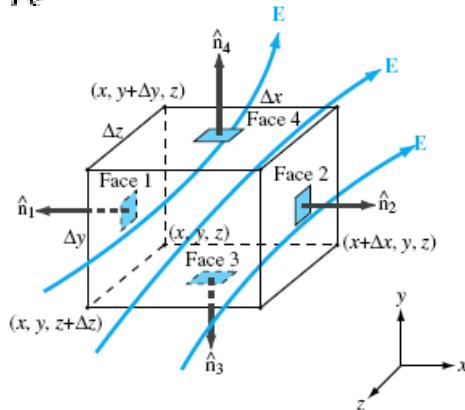
Vector analysis

Divergence

$$\nabla \cdot \mathbf{E} \triangleq \operatorname{div} \mathbf{E} = \frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} + \frac{\partial E_z}{\partial z}$$

$$\operatorname{div} \mathbf{E} \triangleq \lim_{\Delta V \rightarrow 0} \frac{\oint_S \mathbf{E} \cdot d\mathbf{s}}{\Delta V},$$

$$\int_V \nabla \cdot \mathbf{E} dV = \oint_C \mathbf{E} \cdot d\mathbf{s} \quad (\text{divergence theorem}),$$

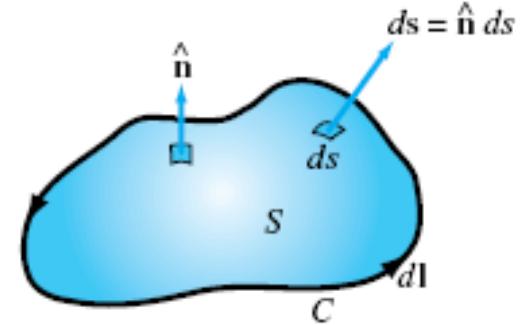


Curl

$$\nabla \times \mathbf{B} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ B_x & B_y & B_z \end{vmatrix}$$

$$\nabla \times \mathbf{B} = \operatorname{curl} \mathbf{B} \triangleq \lim_{\Delta s \rightarrow 0} \frac{1}{\Delta s} \left[\hat{n} \oint_C \mathbf{B} \cdot d\mathbf{l} \right]_{\max}.$$

$$\int_S (\nabla \times \mathbf{B}) \cdot d\mathbf{s} = \oint_C \mathbf{B} \cdot d\mathbf{l} \quad (\text{Stokes's theorem}),$$



Differential Maxwell's equations (1)

Time domain Maxwell's equations:

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} \quad \text{Faraday's law}$$

$$\nabla \times \mathbf{H} = \frac{\partial \mathbf{D}}{\partial t} + \mathbf{J} \quad \text{Ampere's law}$$

$$\nabla \cdot \mathbf{D} = \rho \quad \text{Gauss' law}$$

$$\nabla \cdot \mathbf{B} = 0 \quad \text{Gauss' law-magnetic}$$

Basic quantities:	$\mathbf{E}(t, \mathbf{r})$ – electric field	V/m
	$\mathbf{H}(t, \mathbf{r})$ – magnetic field	A/m
	$\mathbf{D}(t, \mathbf{r})$ – electric flux density	C/m^2
	$\mathbf{B}(t, \mathbf{r})$ – magnetic flux density	webers/m^2
	\mathbf{J} – electric current density	A/m^2
	ρ – charge density	C/m^3

Differential Maxwell's equations (2)

Frequency domain Maxwell's equations $\partial/\partial t \rightarrow j\omega$:

$$\nabla \times \tilde{\mathbf{E}} = -j\omega \tilde{\mathbf{B}} \quad \text{Faraday's law}$$

$$\nabla \times \tilde{\mathbf{H}} = j\omega \tilde{\mathbf{D}} + \tilde{\mathbf{J}} \quad \text{Ampere's law}$$

$$\nabla \cdot \tilde{\mathbf{D}} = \tilde{\rho} \quad \text{Gauss's law}$$

$$\nabla \cdot \tilde{\mathbf{B}} = 0 \quad \text{Gauss's law-magnetic}$$

Basic quantities:	$\tilde{\mathbf{E}}(\omega, \mathbf{r})$ – electric field (phasor)	V/m
	$\tilde{\mathbf{H}}(\omega, \mathbf{r})$ – magnetic field (phasor)	A/m
	$\tilde{\mathbf{D}}(\omega, \mathbf{r})$ – electric flux density (phasor)	C/m^2
	$\tilde{\mathbf{B}}(\omega, \mathbf{r})$ – magnetic flux density (phasor)	webers/m^2
	\mathbf{J} – electric current density (phasor)	A/m^2
	ρ – charge density (phasor)	C/m^3

Differential Maxwell's equations (3)

- **Analogy with TL equations** ($\partial/\partial t, \omega \neq 0$)

— **Assume lossless case** $\Rightarrow \mathbf{B} = \mu_0 \mathbf{H}, \mathbf{D} = \epsilon_0 \mathbf{E}$

— **Time domain**

$$\nabla \times \mathbf{E} = -\mu_0 \frac{\partial \mathbf{H}}{\partial t}$$

$$\nabla \times \mathbf{H} = \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}$$

$$\{\mathbf{E}, \mathbf{H}\} = \underbrace{\{v, i\}(t, z)}_{\text{voltage}} \underbrace{\{\psi_v, \psi_i\}(\mathbf{r})}_{\text{cross-sectional distribution}}$$

$$-\frac{\partial v(t, z)}{\partial z} = L' \frac{\partial i(t, z)}{\partial t}$$

$$-\frac{\partial i(t, z)}{\partial z} = C' \frac{\partial v(t, z)}{\partial t}$$

— **Frequency domain**

$$\nabla \times \tilde{\mathbf{E}} = -j\omega \tilde{\mathbf{B}}$$

$$\nabla \times \tilde{\mathbf{H}} = j\omega \tilde{\mathbf{D}} + \tilde{\mathbf{J}}$$

$$\{\tilde{\mathbf{E}}, \tilde{\mathbf{H}}\} = \underbrace{\{\tilde{V}, \tilde{I}\}(\omega, z)}_{\text{voltage}} \underbrace{\{\psi_v, \psi_i\}(\mathbf{r})}_{\text{cross-sectional distribution}}$$

$$-\frac{d\tilde{V}(\omega, z)}{dz} = j\omega L' \tilde{I}(\omega, z)$$

$$-\frac{d\tilde{I}(\omega, z)}{dz} = j\omega L' \tilde{V}(\omega, z)$$

Differential Maxwell's equations (4)

- **Statics**
 - $\partial/\partial t = \omega = 0 \Rightarrow$ two uncoupled sets of (static) equations are obtained for \mathbf{E} & \mathbf{D} and \mathbf{H} & \mathbf{B}
 - **Electrostatics**

$$\nabla \times \mathbf{E} = 0$$

$$\nabla \cdot \mathbf{D} = \rho$$

- **Magnetostatics**

$$\nabla \times \mathbf{H} = \mathbf{J}$$

$$\nabla \cdot \mathbf{B} = 0$$

Integral Maxwell's equations (1)

Time domain Maxwell's equations:

$$\oint_C \mathbf{E} \cdot d\mathbf{l} = -\frac{d}{dt} \iint_S \mathbf{B} \cdot d\mathbf{s}$$

Faraday's law (Stock's theorem to FD)

$$\oint_C \mathbf{H} \cdot d\mathbf{l} = \frac{d}{dt} \iint_S \mathbf{D} \cdot d\mathbf{s} + I$$

Ampere's law (Stock's theorem to FD)

$$\iint_S \mathbf{D} \cdot d\mathbf{s} = Q$$

Gauss's law (Gauss' theorem to FD)

$$\iint_S \mathbf{B} \cdot d\mathbf{s} = 0$$

Gauss's law-magnetic (Gauss' theorem to FD)

Basic quantities:

$I(t)$ – total current A

$Q(t)$ – total charge C

Integral Maxwell's equations (2)

Frequency domain Maxwell's equations:

$$\oint_C \tilde{\mathbf{E}} \cdot d\mathbf{l} = -j\omega \iint_S \tilde{\mathbf{B}} \cdot d\mathbf{s}$$

Faraday's law (Stock's theorem to FD)

$$\oint_C \tilde{\mathbf{H}} \cdot d\mathbf{l} = j\omega \iint_S \tilde{\mathbf{D}} \cdot d\mathbf{s} + \tilde{I}$$

Ampere's law (Stock's theorem to FD)

$$\iint_S \tilde{\mathbf{D}} \cdot d\mathbf{s} = \tilde{Q}$$

Gauss' law (Gauss' theorem to FD)

$$\iint_S \tilde{\mathbf{B}} \cdot d\mathbf{s} = 0$$

Gauss' law-magnetic (Gauss' theorem to FD)

Basic quantities:

$\tilde{I}(\omega)$ – total current A

$\tilde{Q}(\omega)$ – total charge C

Integral Maxwell's equations (3)

- **Statics**
 - **Electrostatics**

$$\oint_C \tilde{\mathbf{E}} \cdot d\mathbf{l} = 0$$

$$\iint_S \tilde{\mathbf{D}} \cdot d\mathbf{s} = \tilde{Q}$$

- **Magnetostatics**

$$\oint_C \tilde{\mathbf{H}} \cdot d\mathbf{l} = \tilde{I}$$

$$\iint_S \tilde{\mathbf{B}} \cdot d\mathbf{s} = 0$$

Charges

- Charge densities

- Volume charge density ρ_v

- $\rho_v = \lim_{\Delta v \rightarrow 0} \frac{\Delta q}{\Delta v} = \frac{dq}{dv}$ (C/m³),

- Average charge per (physically large) unit volume

- Surface charge density ρ_s

- $\rho_s = \lim_{\Delta s \rightarrow 0} \frac{\Delta q}{\Delta s} = \frac{dq}{ds}$ (C/m²),

- Average charge per unit surface

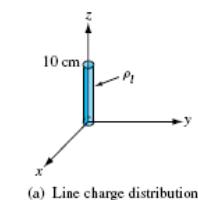
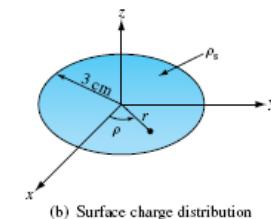
- Line charge density ρ_l

- $\rho_l = \lim_{\Delta l \rightarrow 0} \frac{\Delta q}{\Delta l} = \frac{dq}{dl}$ (C/m)

- Average charge per unit length

- Total charge

$$Q = \int_v \rho_v dv \text{ or/and } Q = \int_s \rho_s ds \text{ or/and } Q = \int_l \rho_l dl$$



Currents

- Current densities
 - Volume charge density

$$\Delta q' = \rho_v \Delta v = \rho_v \Delta l \Delta s' = \rho_v u \Delta s' \Delta t,$$

- Charge through a generalized cross-section

$$\Delta q = \rho_v \mathbf{u} \cdot \Delta \mathbf{s} \Delta t,$$

- Current density

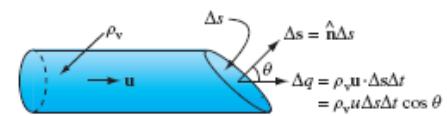
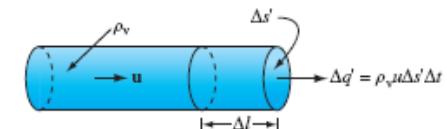
$$\Delta I = \frac{\Delta q}{\Delta t} = \rho_v \mathbf{u} \cdot \Delta \mathbf{s} = \mathbf{J} \cdot \Delta \mathbf{s},$$

$$\mathbf{J} = \rho_v \mathbf{u} \quad (\text{A/m}^2)$$

- Surface current density

- thin current sheet (thickness $d \rightarrow 0$) $\Rightarrow \mathbf{J}_s = \lim_{d \rightarrow 0} d\mathbf{J}$ (A/m)

- Total current $I = \int_S \mathbf{J} \cdot \hat{\mathbf{n}} ds$ or/and $I = \int_l \mathbf{J}_s \cdot \hat{\mathbf{n}} dl$



Vector wave equation

Lossless case: $\mathbf{D} = \epsilon_0 \mathbf{E}$, $\mathbf{B} = \mu_0 \mathbf{H}$, $\rho = 0$, $\mathbf{J} = 0$

$$\nabla \times \mathbf{E} = -\mu_0 \frac{\partial \mathbf{H}}{\partial t} \quad \longrightarrow \quad \nabla \times (\nabla \times \mathbf{E}) = -\frac{\partial}{\partial t} (\nabla \times \mathbf{H})$$

use $\nabla \times (\nabla \times \mathbf{E}) = \nabla (\nabla \cdot \mathbf{E}) - \nabla^2 \mathbf{E} = -\nabla^2 \mathbf{E}$

0 ↗

$$-\nabla^2 = -\mu_0 \frac{\partial}{\partial t} (\nabla \times \mathbf{H}) \qquad \text{use} \qquad \nabla \times \mathbf{H} = \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}$$

$$\nabla^2 \mathbf{E} - \mu_0 \epsilon_0 \frac{\partial^2 \mathbf{E}}{\partial t^2} = 0 \xrightarrow{\text{for 1D}} \frac{\partial^2 \mathbf{E}}{\partial z^2} - \frac{1}{c^2} \frac{\partial^2 \mathbf{E}}{\partial t^2} = 0$$

Electric field boundary conditions (1)

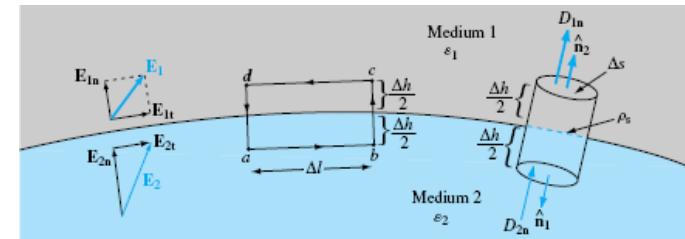
- Boundary conditions for any dissimilar media

- Tangential components

$$\oint_C \mathbf{E} \cdot d\mathbf{l} = -d/dt \iint_S \mathbf{B} \cdot d\mathbf{s} \quad \Delta h \rightarrow 0,$$

$$\Rightarrow \int_a^b \mathbf{E}_2 \cdot d\mathbf{l} + \int_c^d \mathbf{E}_1 \cdot d\mathbf{l} + \underbrace{\int_a^c \mathbf{E} \cdot d\mathbf{l}}_{\sim \Delta h \rightarrow 0} + \underbrace{\int_c^b \mathbf{E} \cdot d\mathbf{l}}_{\sim \Delta h \rightarrow 0} = -\frac{d}{dt} \iint_S \mathbf{B} \cdot d\mathbf{s}$$

$$E_{2t} \Delta l - E_{1t} \Delta l = 0 \Rightarrow E_{1t} = E_{2t} \Rightarrow \hat{\mathbf{n}}_2 \times (\mathbf{E}_1 - \mathbf{E}_2) = 0$$



- Normal components

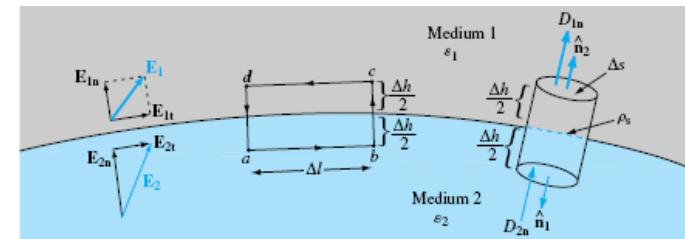
$$\oint_S \mathbf{D} \cdot d\mathbf{s} = \int_{\text{top}} \mathbf{D}_1 \cdot \hat{\mathbf{n}}_2 ds + \int_{\text{bottom}} \mathbf{D}_2 \cdot \hat{\mathbf{n}}_1 ds = Q = \rho_s \Delta s$$

$$\hat{\mathbf{n}}_2 \cdot (\mathbf{D}_1 - \mathbf{D}_2) = \rho_s$$

Electric field boundary conditions (2)

- **Boundary conditions for dielectric media**
 - Tangential components

$$E_{1t} = E_{2t} \Rightarrow \hat{\mathbf{n}}_2 \times (\mathbf{E}_1 - \mathbf{E}_2) = 0$$



- Normal components

$$\hat{\mathbf{n}}_2 \cdot (\mathbf{D}_1 - \mathbf{D}_2) = 0$$

Electric field boundary conditions (3)

- **Dielectric-conductor boundary conditions**
 - Assume medium 2 is perfect conductor

$$\Rightarrow \mathbf{E}_2 = \mathbf{D}_2 = 0,$$

$$\Rightarrow \begin{aligned} E_{1t} &= D_{1t} = 0, \\ D_{1n} &= \epsilon_1 E_{1n} = \rho_s. \end{aligned}$$

$$\Rightarrow \boxed{\begin{aligned} \hat{\mathbf{n}}_2 \times \mathbf{E}_1 &= 0 \\ \hat{\mathbf{n}}_2 \cdot \mathbf{D}_1 &= \rho_s \end{aligned}}$$

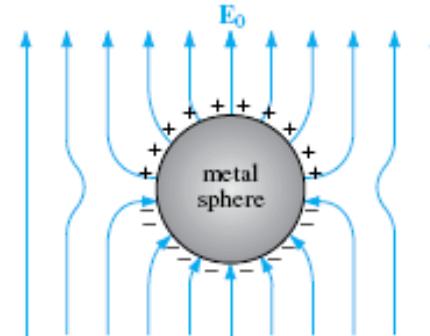


Figure 4-21: Metal sphere placed in an external electric field \mathbf{E}_0 .

Magnetic field boundary conditions

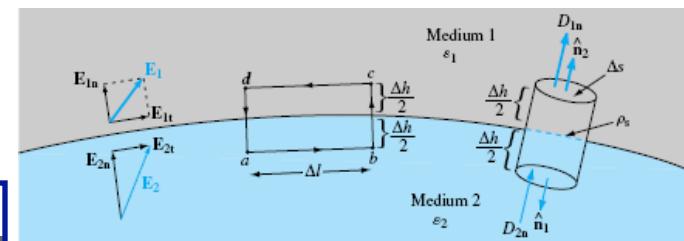
- **Boundary conditions for any dissimilar media**
 - Tangential components

$$\oint_C \mathbf{H} \cdot d\mathbf{l} = \int_a^b \mathbf{H}_2 \cdot d\mathbf{l} + \int_c^d \mathbf{H}_1 \cdot d\mathbf{l} = I,$$

$$H_{2t} \Delta l - H_{1t} \Delta l = J_s \Delta l$$

$$H_{2t} - H_{1t} = J_s$$

$$\hat{\mathbf{n}}_2 \times (\mathbf{H}_1 - \mathbf{H}_2) = \mathbf{J}_s$$



- Normal components

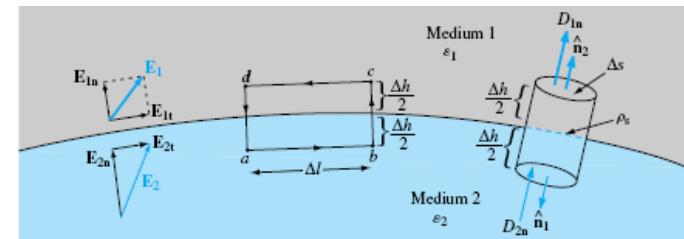
$$\oint_S \mathbf{B} \cdot d\mathbf{s} = 0$$

$$\hat{\mathbf{n}}_2 \cdot (\mathbf{B}_1 - \mathbf{B}_2) = 0$$

Magnetic field boundary conditions

- **Boundary conditions for dielectric media**
 - Tangential components

$$\hat{\mathbf{n}}_2 \times (\mathbf{H}_1 - \mathbf{H}_2) = 0$$



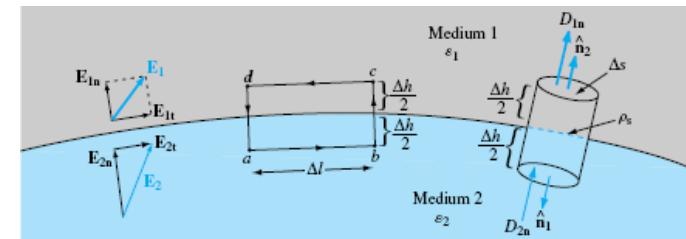
- Normal components

$$\hat{\mathbf{n}}_2 \cdot (\mathbf{B}_1 - \mathbf{B}_2) = 0$$

Magnetic field boundary conditions

- **Boundary conditions for perfect conductors**
 - Tangential components

$$\hat{\mathbf{n}}_2 \times \mathbf{H}_1 = \mathbf{J}_S$$



- Normal components

$$\hat{\mathbf{n}}_2 \cdot \mathbf{B}_1 = 0$$

Complete set of differential MEs

Maxwell's equations:

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} \quad \text{Faraday's law}$$

$$\nabla \times \mathbf{H} = \frac{\partial \mathbf{D}}{\partial t} + \mathbf{J} \quad \text{Ampere's law}$$

$$\nabla \cdot \mathbf{D} = \rho \quad \text{Gauss' law}$$

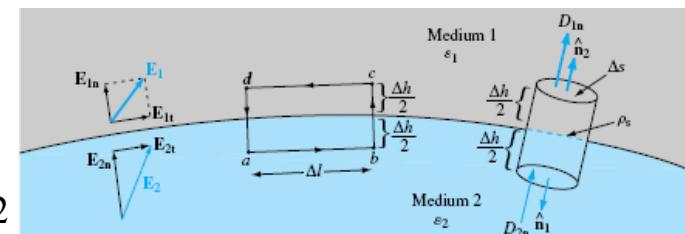
$$\nabla \cdot \mathbf{B} = 0 \quad \text{Gauss' law-magnetic}$$

Boundary conditions:

$$\hat{\mathbf{n}}_{12} \times (\mathbf{E}_2 - \mathbf{E}_1) = 0 \quad \hat{\mathbf{n}}_{12} \times (\mathbf{H}_2 - \mathbf{H}_1) = \mathbf{J}_s$$

$$\hat{\mathbf{n}}_{12} \cdot (\mathbf{D}_2 - \mathbf{D}_1) = \rho_s \quad \hat{\mathbf{n}}_{12} \cdot (\mathbf{B}_2 - \mathbf{B}_1) = 0$$

$\hat{\mathbf{n}}_{12} = \hat{\mathbf{n}}_2$ — normal to the interface from medium 1 to medium 2



Constitutive relations: $\mathbf{D} = \epsilon \mathbf{E}$

$$\mathbf{B} = \mu \mathbf{H}$$