

1 On the Dynamics of Global Temperature

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4 **Abstract**

5 In this alternative theory of global temperature dynamics over the
6 annual to the glacial time scales, the accumulation of variations in so-
7 lar irradiance dominates the dynamics of global temperature change. A
8 straightforward recurrence matrix representation of the atmosphere/surface/deep
9 ocean system, models temperature changes by (1) the size of a forc-
10 ing, (2) its duration (due to accumulation of heat), and (3) the depth
11 of forcing in the atmosphere/surface/deep ocean system (due to in-
12 creasing mixing losses and increasing intrinsic gain with depth). The
13 model can explain most of the rise in temperature since 1950, and more
14 than 70% of the variance with correct phase shift of the 11-year solar
15 cycle. Global temperature displays the characteristics of an accumula-
16 tive system over 6 temporal orders of magnitude, as shown by a linear
17 f^{-1} log-log relationship of frequency to the temperature range, and
18 other statistical relationships such as near random-walk and distribu-
19 tion asymmetry. Over the last century, annual global surface tempera-

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20 ture rises or falls $0.063 \pm 0.028C/W/m^2$ per year when solar irradiance
21 is greater or less than an equilibrium value of $1366W/m^2$ at top-of-
22 atmosphere. Due to an extremely slow characteristic time scale the
23 notion of 'equilibrium climate sensitivity' is largely superfluous. The
24 theory does not require a range of distinctive feedback and lag param-
25 eters. Mixing losses attenuate the effectiveness of greenhouse gasses.
26 Most recent warming can be explained without recourse to increases
27 in heat-trapping gases produced by human activities.

28 **1 Introduction**

29 An accumulating body of evidence showing high sensitivity of global tem-
30 perature to solar variations, accounting for more than half of global warm-
31 ing since the mid 20th century Douglass and Clader [2002], Shaviv [2008],
32 Scafetta and West [2007], Scafetta [2009, 2010a,b], and the recent (2003
33 – 2011) flat warming/cooling rate of $0.1 \pm 0.2W/m^2$ ocean heat content
34 anomaly instead of the frequently-cited, large, positive computed radiative
35 imbalance of $0.6 \pm 0.3W/m^2$ Knox and Douglass [2010], demands an expla-
36 nation. The absence of a known mechanism to explain why Earth's temper-
37 ature may be 'mysteriously hypersensitive to solar variations' Scafetta et al.
38 [2009], Duffy et al. [2009] blocks the acceptance of these empirical results
39 Lockwood and Fröhlich [2008].

40 Another related problem is the range of estimates of climate sensitivity
41 between climate models IPCC [2007] and natural experiments Idso [1998],
42 natural cycles Scafetta [2010b], Douglass and Clader [2002], the Pinatubo

43 event Douglass and Knox [2005], Douglass et al. [2006b], Bender et al. [2010],
44 and temperature fluctuations Stott et al. [2003], Lindzen and Choi [2009],
45 Spencer and Braswell [2010], Dessler [2010]. These issues, together with the
46 difficulty of reducing the range of estimates of the sensitivity of the system
47 to CO_2 doubling to less than an order of magnitude IPCC [2007] suggests a
48 fundamental gap in our understanding of climate sensitivity.

49 Global temperature variations are conventionally thought of as the in-
50 teraction of different types of forcings and feedbacks over different time
51 scales: fast, medium and slow depending on type, water vapour, GHGs,
52 and ice-sheet albedo Hansen et al. [2011]. Model sensitivity ranges from
53 $0.75C/W/m^2$ to $8C/W/m^2$ and greater for long term equilibria Stern [2005].
54 General Circulation Models (GCMs) attempt to represent the complexity of
55 these physical relationships and values, both measured and presumed.

56 Natural and modeled systems contain a mix of fast and slow equilibrating
57 components. They have a crucial difference. If fast, then continued forcing
58 at the same average level does not cause any additional warming; forcing is
59 directly related to response. If slow, constant high levels can cause ongoing
60 warming until equilibrium is reached. In the slow case, the forcing cannot be
61 directly related to the response, and requires a slightly more complex model,
62 here an autocorrelated (AR) recurrence equation or matrix. We claim here
63 (1) the slow, equilibrium component is largely due to the simple, physical,
64 accumulation of heat; (2) the accumulative mechanism is responsible for high
65 sensitivity to solar radiation, and consequently; (3) insufficient consideration
66 of slow dynamics has lead to errors and underestimation of the contribution

67 of solar forcing to climate change.

68 Energy Balance Models (EBMs) are simple models of the temperature
69 response of a body to energy coming in and going out, in this case the Earth
70 Knox and Douglass [2010]. We will consider the conjecture (called the 'Ac-
71 cumulation Theory of Climate Change') that accumulated energy provides
72 a parsimonious and physically motivated explanation for a wide range of
73 climate observations. The development of the theory draws on notions from
74 control theory Stubberud et al. [1994] and the electronic integration am-
75 plifier. In the first section, we show evidence the theory explains, and the
76 physical representation of near-random walk behavior of the Earth's tem-
77 perature. We then explain the inverse relationship of temperature variance
78 to the time scale (f^{-1}) over at least six orders of magnitude.

79 This approach is also justified by cointegration theory, wherein variables
80 can only be related if they have the same long-run behaviour, i.e they are
81 both stable (stationary or zero trend) or both unstable (non-stationary or
82 trending). Integration order $I(n)$ is the number of differentiations n required
83 to make the variable stable. It has been shown that while solar intensity is
84 $I(0)$, global temperature requires one differentiation to become stable, and
85 is, therefore, $I(1)$ Beenstock and Reingewertz [2010]. Order difference is like
86 apples and oranges; the correct approach is to bring them to the same order
87 by integrating solar intensity before regressing with global temperature, as
88 we do here.

89 At the contemporary time scale, solar irradiance since 1950 has been
90 above average Usoskin et al. [2003], Solanki et al. [2004]. We show that

91 accumulation of the above average solar isolation can explain the majority of
92 the magnitude and phase of temperature variations from 1950 to the present.
93 The AR coefficient of atmospheric temperature increases with decreasing
94 height, increasing linearly with increasing density, from 0.2 to 0.5 in the
95 upper and lower troposphere, to 0.9 at the surface and more in the deep
96 ocean. The graduated AR system creates multiple characteristic decay times
97 represented by a recurrence matrix model.

98 The use of a matrix recurrence formula to analyse climate sensitivity
99 in the climate system is a new approach, growing out of previous work on
100 stochastic models and the spectral scaling of climate variability Vjushin et al.
101 [2002], Stockwell [2006], Koutsoyiannis and Cohn [2008], and empirical stud-
102 ies of derivatives and integrals of tropical ocean indices and multi-decadal
103 warming trends McLean et al. [2009], Stockwell and Cox [2009].

104 **2 Observations**

105 **2.1 Paleoclimatic Time Scale**

106 A good theory should explain all the available observations with as few
107 parameters as possible. The Earth's temperature record spans over 10^6 to
108 10^{-2} years. Global temperature data sets ranging over the 30 years monthly
109 variations of satellite records, to 150 year surface temperature records, 1000
110 year proxy climate records, and the 800,000 years EPICA ice core record.
111 The variance or range or spectral power of these data varies with frequency
112 f (or wavelength f^{-1}) (Fig. 1).

113 The time series needs some preparation due to spectral biases. The
114 EPICA record was aggregated to 1000-year means and dividing by 4 to
115 account for polar bias. The proxy records Loehle [2007], Moberg et al.
116 [2005] were aggregated to 20 year means.

117 A spectral plot of the data sets (Fig 1) shows a single, linear relationship
118 (solid grey line) with decline of almost f^{-1} (dashed grey line). Remarkably,
119 the standard deviation (SD) of temperature over 6 orders of magnitude
120 appears to be proportional to the log of the wavelength:

$$121 \quad SD(t) \propto \log(\text{wavelength}).$$

122 For example, if the standard deviation (σ) is 0.1C at the scale of one year,
123 then natural variation over 10 years will be 0.23C, 100 years will be 0.46C,
124 1000 years will be 0.69C and over the million years, 1.38C. These figures are
125 consistent with a range of around 5C for glacial-interglacial transitions. The
126 power spectrum may deviate from the linear relationship at the maximum
127 and minimum ends as shown.

128 Low total variation in solar irradiation (TSI) is available to generate
129 such such temperature variations. TSI is around $1366W/m^2$ at the top of
130 atmosphere, but the geometry of the globe and an average albedo of 30%
131 reduces surface TSI to 17% or $235W/m^2$. Variations in TSI over the 11 year
132 solar cycle typically do not exceed 0.1% or $0.24W/m^2$ at the surface. The
133 increase in TSI since the last century is about $0.28W/m^2$, while variation
134 in TSI varies no more than 0.2% or $0.5W/m^2$ at the surface over 100,000
135 years due to orbital variations Lean and Rind [2001], Muscheler et al. [2007].
136 Geological forcing is of smaller magnitude, estimated at $0.09W/m^2$.

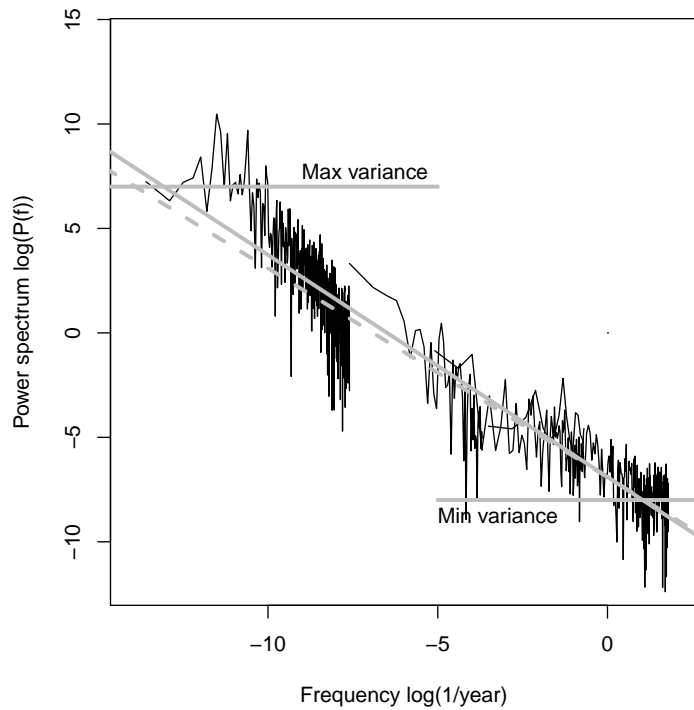


Figure 1: Power spectrum of estimates of global temperature ranging from 800K years (EPICA ice core) to the 30 years (satellite record). The line of best fit through the data (gray) does not differ significantly from the theoretically expected inverse of frequency f^{-1} found in integration amplifiers. Horizontal gray lines indicate possible maximum and minimum amplitudes at approximately 22,000 years and 1 year.

137 The temperature sensitivity of a black body to forcing is $0.25C/W/m^2$
138 to $0.3C/W/m^2$. Therefore, the direct, proportional effect of an increase in
139 solar forcing of $0.2W/m^2$ would be only $0.05C$. Conventionally, small changes
140 in temperature launch feedbacks from water vapor, greenhouse gases and
141 surface albedo that promote the response of the Earth's atmosphere by an
142 order of magnitude.

143 In the accumulative theory, accumulation of heat in the mass of the land
144 and ocean causes global temperature increases. Both the magnitude and the
145 duration of a forcing determines the heat accumulated and, therefore, the
146 temperature.

147 A 'back of envelope' calculation indicates that a solar forcing of $0.1W/m^2$
148 for 1 year transfers 3.1×10^6 Joules of heat (3.1×10^6 sec in a year) to the ocean.
149 Based on the specific heat of water (4.2 J/gK) and the number of grams in a
150 cubic meter of water (10^6), a water column 100 m deep would warm $0.008K$
151 in one year, or $0.8K$ in a century. Thus, the 20th century temperature rise
152 can be explained by the accumulation of an above average solar forcing of
153 $0.1W/m^2$ in the ocean over the period Lean and Rind [2001, 2008].

154 Similarly, a forcing of $0.1W/m^2$ accumulated for 5,000 years would in-
155 crease the whole ocean temperature by $4C$, sufficient to account for inter-
156 glacial warming.

157 Control systems theory provides an elegant and powerful formalism for
158 describing such systems. The Bode plot (Fig. 2) illustrates the spectral
159 energy (upper), and phase relationship (lower) to the frequency (or wave-
160 length) that simultaneously provides a complete picture of the way a system

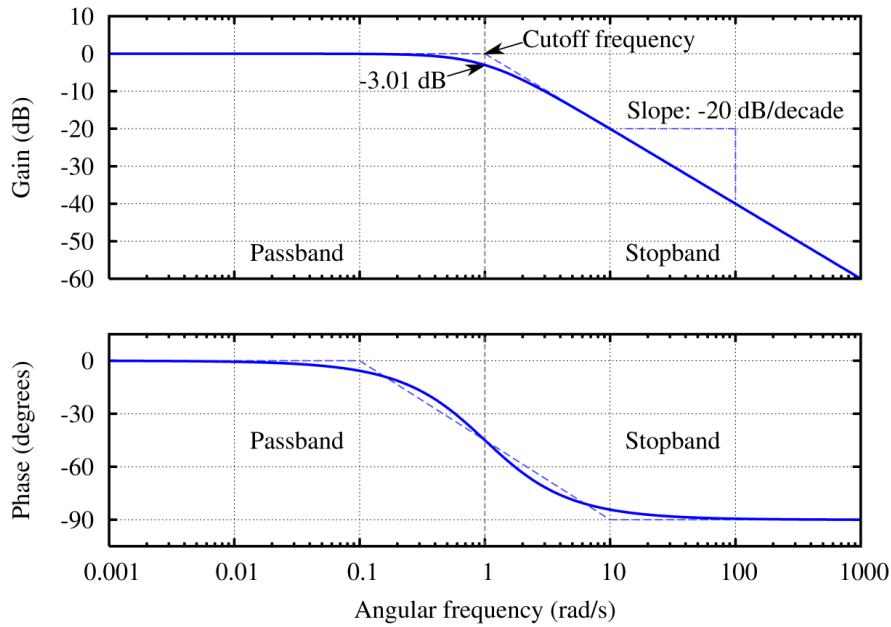


Figure 2: A Bode plot for an integrative amplifier, or low-pass filter, showing the amplification increases linearly with decreasing frequency until it reaches an amplification limit at a cutoff frequency. The phase also varies with frequency (from Wiki commons).

161 modifies input Stubberud et al. [1994]. The frequency response of a ideal
 162 integrator is a downward-sloping line, indicating higher gain at lower fre-
 163 quencies:

164 T1: f^{-1}

165 A finite accumulator has a plateau of maximum gain (Fig 2). Basic con-
 166 trol theory shows how the function below can be obtained from the Laplace
 167 transform of the basic energy balance model (4) described in the next section
 168 Stubberud et al. [1994].

169 T2: $(\alpha + f)^{-1}$

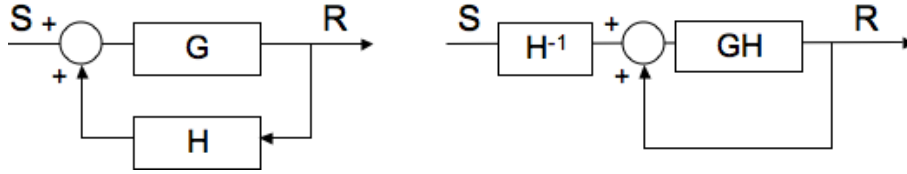


Figure 3: (A) A feedback controller with system components G and feedback H , and (B) the corresponding diagram rearranged into an integration amplifier, where the feedback is a unit loop.

170 T1 is a ideal integrator with an infinite gain at zero (low) frequency. T2
 171 has finite gain via the inclusion of a small system loss. 'Reddening' is a bias
 172 of random noise towards low frequencies Roe [2009]. In comparison, random
 173 or white noise has equal power at all frequencies, producing a horizontal line
 174 on the spectral plot.

175 We can now explain the maximum and minimum limits to the integrator
 176 in the spectral plot of global temperature (Fig 1). The maximum response
 177 appears at about period e^{10} or 22,000 years, suggesting the limits of the
 178 temperature range. The minimum variance appears at time periods less
 179 than one year suggestive of white noise. The Earth system appears to act
 180 as an ideal integration amplifier between these time scales.

181 Control systems theory formalizes the notion of feedbacks (Fig. 3). The
 182 simplest loop has two transfer functions, a transient system amplification G
 183 and a feedback H summed at the control point (A). H and G are simply con-
 184 stants. The recurrence equation $Y_{i+1} = G(S + HY_i)$ describes one iteration
 185 around the loop:

186
$$Y_{i+1} = GHY_i + GS$$

187 An integration amplifier, the block diagram in Fig 3 (A), can be rear-
188 ranged as a unit loop (B), summing output and input. Though this rear-
189 rangement, amplification is a single transfer function $\frac{1}{H} * GH$ or G embedded
190 in an iterative sum. The two views are equivalent; feedback can be trans-
191 formed into an integrating amplifier.

192 The example demonstrates that integration is more specific than feed-
193 back; asserting the dominant mechanism of climate change is an accumula-
194 tion mechanism is a stronger claim than invoking feedbacks. For example, in
195 classical feedback an increase in isolation of $0.1W/m^2$ produces more water
196 vapour in the air which permits absorption of $0.2W/m^2$ of radiation. In the
197 accumulative view, increasing humidity and temperature can be regarded as
198 the accumulated stock of heat.

199 **2.2 Recent Warming**

200 The conventional view of global warming is that there is no plausible expla-
201 nation for the rise in temperature since 1950 other than the heat-trapping
202 effects of human emissions of greenhouse gasses IPCC [2007]. We explore a
203 model of accumulated solar radiation denoted $\Sigma Solar$ (or CumSolar on the
204 Figures) Lean [2001], also a sunspot series due to uncertainty in the solar
205 irradiance satellite composites Scafetta [2009].

206 Exhaustive search for the equilibrium value was optimized on the cor-
207 relation of accumulated monthly irradiance (and sunspots) with HadCRU
208 temperature from 1950 and to present. The zero point was $1365.9 W/m^2$
209 and 21 monthly sunspots; temperature is generally rising above and gener-

210 ally falling below these values. The equilibrium value was subtracted from
211 the solar irradiance series before accumulation, producing $\Sigma Solar$.

212 Figures 4, 5, 6, and 7 show the linear regressions of $\Sigma Solar$ against
213 global temperature datasets. The $\Sigma Solar$ variable is highly significant and
214 accounts for more than 60% (blue) and 70% (red including volcanics) of the
215 variation (Fig. 7). By comparison, the direct irradiance has little correlation
216 ($R^2 < 0.1$) (orange). Solar sunspot counts gave similar results.

217 An additional regression including a time term provided the opportunity
218 for another trending factor, such as increasing concentrations of greenhouse
219 gasses, to explain the trend. If the trend term was significant, then the
220 relative contribution of $\Sigma Solar$ to the warming trend would be estimated
221 by subtracting the residual trend from the overall trend. The residual trend
222 term was not significant in any of the data sets. The accumulated radiance
223 was $128 W/m^2$ during the 58 years since 1950, representing an average solar
224 forcing of $0.18 W/m^2$ at the top of the atmosphere (TOA) (i.e. $128/58$
225 years/12 months).

226 The $\Sigma Solar$ variable also has the correct phase relationship with tem-
227 perature, as shown by the peak of the cross-correlation at zero lag (Fig. 8).
228 The volcanics and direct sun spots have low, lagged correlation. Changes
229 in ocean temperatures lead land temperatures by one year, indicative of an
230 ocean-controlled system.

231 Fig. 9 illustrates the main assumed atmospheric forcings from the GISS
232 general circulation model (GCM) (as listed in the file RadF.txt) Hansen
233 et al. [2011] and brought up to date from 2003: well mixed green house

GISS Global Temperature and Accumulated Solar Radiation

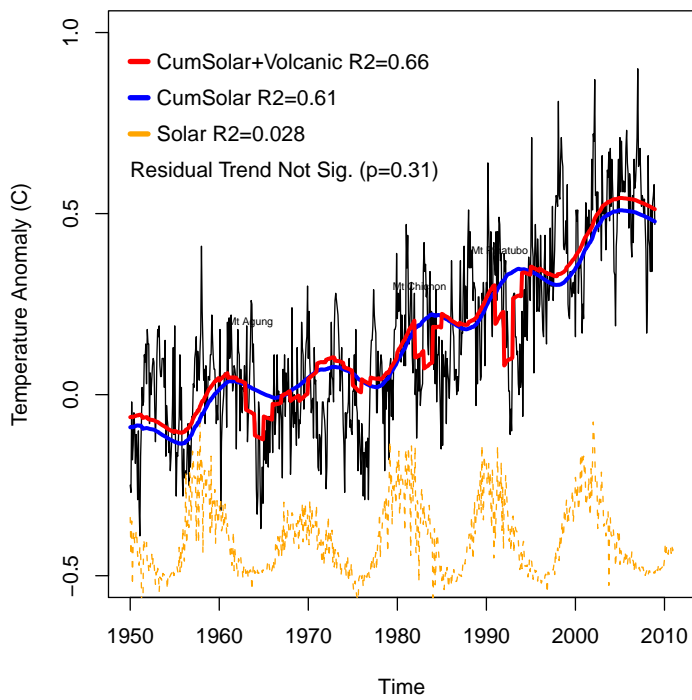


Figure 4: Cumulative solar irradiance (blue) and volcanic forcing (red) is highly correlated with GISS sea surface global temperature and explains the trend in temperature since 1950. The direct solar irradiance (orange) is uncorrelated with temperature.

CRU Sea Global Temperature and Accumulated Solar Radiation

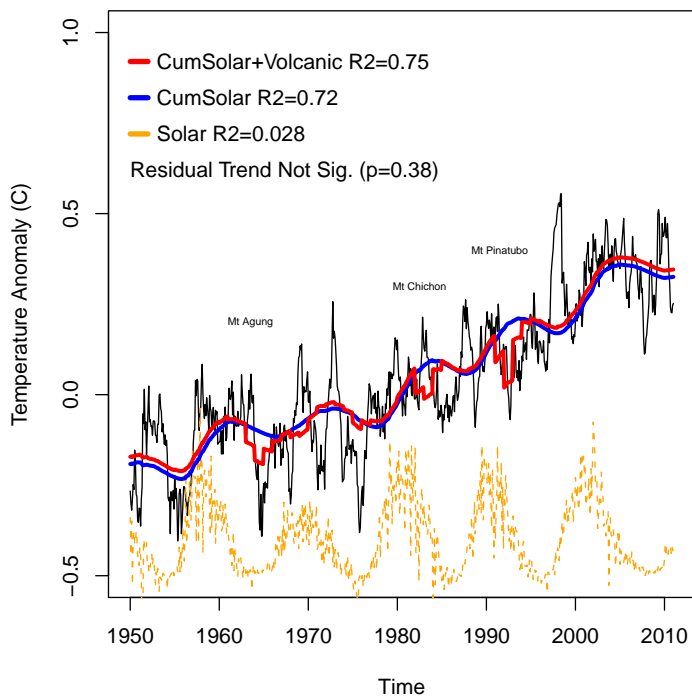


Figure 5: Cumulative solar irradiance (blue) and volcanic forcing (red) is highly correlated with HadSST sea surface global temperature and explains the trend in temperature since 1950. The direct solar irradiance (orange) is uncorrelated with temperature.

CRU Land Global Temperature and Accumulated Solar Radiation

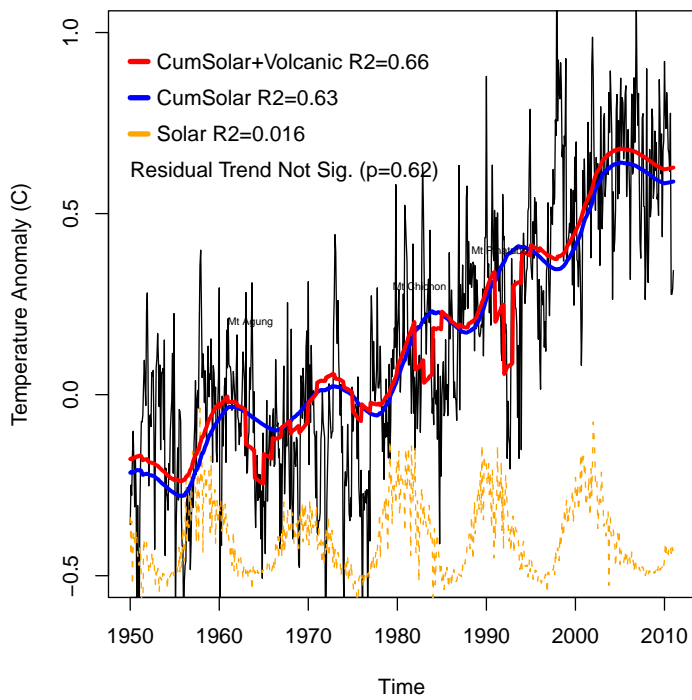


Figure 6: Cumulative solar irradiance (blue) and volcanic forcing (red) is highly correlated with HadLST land surface global temperature and explains the trend in temperature since 1950. The direct solar irradiance (orange) is uncorrelated with temperature.

HadCRU Global Temperature and Accumulated Solar Radiation

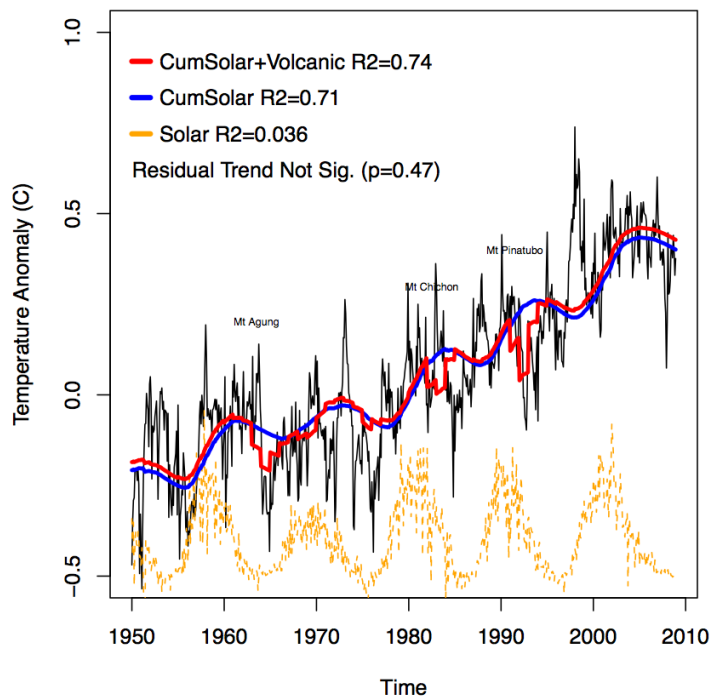


Figure 7: Cumulative solar irradiance (blue) and volcanic forcing (red) is highly correlated with HadCRU global temperature and explains the trend in temperature since 1950. The direct solar irradiance (orange) is uncorrelated with temperature.

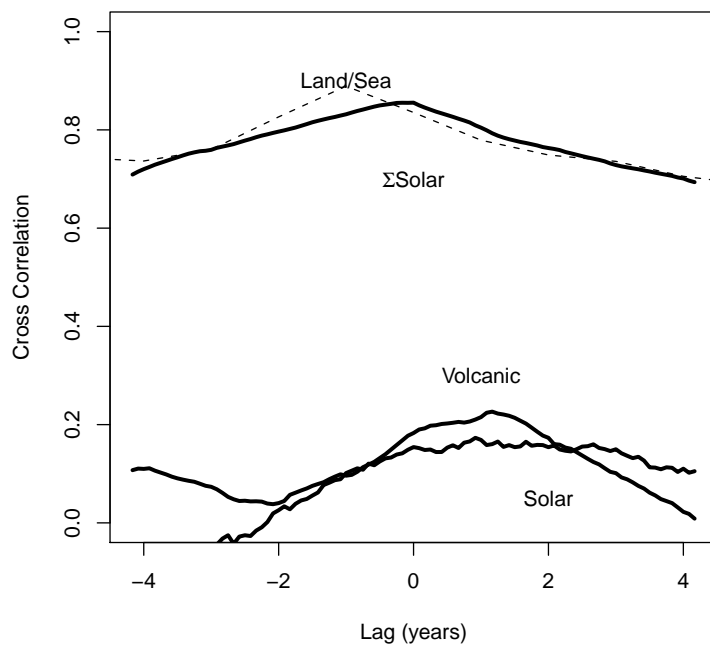


Figure 8: The cross correlation relationships showing the high correlation and correct phase of the accumulated solar radiance. The volcanics are lower correlation and lagged while the direct solar irradiance has very low, lagged correlation. Also shown is the one year lead of ocean over land temperatures (dashed).

234 gasses, W.M_GHG_s, which warmed the Earth about 1.3C, stratospheric
235 aerosols (volcanic) StratAer, and reflective aerosols ReflAer which cooled
236 it about 0.7C. Regression without $\Sigma Solar$ replicates the assumed contribu-
237 tion of W.M_GHG_s and aerosols in the GISS dataset (dashed lines in Fig.9).

238 The contribution of W.M_GHG_s drops to less than half with $\Sigma Solar$ in
239 the regression (red arrow). Thus, a combination of forcings with lower esti-
240 mates of CO_2 sensitivity and higher solar contribution would be consistent
241 with both the rise in global temperature, and the flattening of temperatures
242 and ocean heat uptake in the last decade Loehle [2009], Douglass and Knox
243 [2009], Knox and Douglass [2010].

244 **2.3 Climate cycles and other effects**

245 The most prominent climate cycles are the solar (Schwabe) cycle averaging
246 11 years, the Pacific Decadal Oscillation (PDO) and the Atlantic Merid-
247 ional Oscillation (AMO) with a quasi-periodicity of around 60 years, and
248 Milankovitch cycles of period around 100,000 years related to the transition
249 between ice ages.

250 The accumulation theory suggests that the amplitude of the cycles should
251 be directly related to their duration, as appears to be the case (Fig 1). Con-
252 trol theory also predicts peak power at $2\pi\tau$ or about six times the charac-
253 teristic decay time Stubberud et al. [1994] related to the 11 year solar cycle
254 and the 60 year PDO and AMO oscillations.

255 Further, control theory shows that the start of the maximum plateau

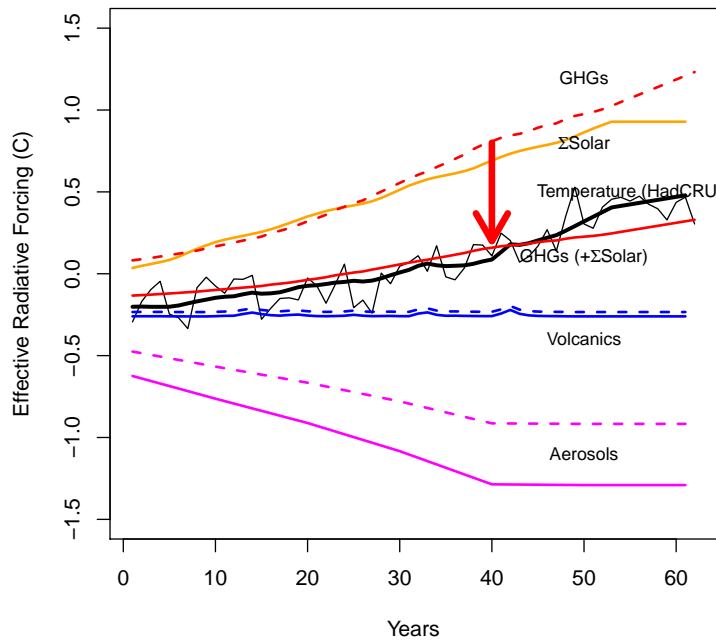


Figure 9: Climate forcings and their contributions to temperature change from the GISS forcing estimates. Regression of GHGs (red) aerosols (purple) and stratospheric aerosols (volcanics - blue) with the observed increase in temperature (black and gray) replicates the assumed contributions in Hansen et al. [2011]. Inclusion of cumulative solar (orange) in the regression decreases GHG contribution by more than half and increases the aerosol contribution.

256 region in the spectral plot in Fig. 1 is at 2π times the characteristic decay
257 time. Thus, the dominant decay time of the system is around 3500 years,
258 which would give an AR coefficient of $a = 0.99971$ – extremely close to a
259 random walk but still stationary (stable).

260 Further evidence in of the model is controlled output and uncontrolled
261 input, as seen in assymmetric temperature changes. Asymmetry is indicated
262 when the mean of differences exceeds the median of the differences. The
263 mean of the EPICA data is 0.01C, greater than the median of -0.07C. In
264 HadCRUT, the mean is 0.005C, greater than the median of 0.003C. In TLT,
265 the mean is 0.002C, greater than the median of -0.009C, indicating asym-
266 metry and hence output control.

267 The accumulation theory does not ignore added forcing, such as interac-
268 tions between solar emanations and the Earth’s magnetic field such as mod-
269 ification of cloud albedo by high-energy particles Svensmark [2007]. Rather,
270 the accumulation theory defines the basic functioning of the system, while
271 indirect solar effects, cloud albedo variations, and aerosols only serve to
272 change the intensity of radiative inputs to the system. Small forcing over
273 long periods may control the timing of cycles.

274 **3 The Models**

275 **3.1 Recurrence Models**

276 A random walk is a mathematical formalisation of trajectories that accu-
277 mulate successive random shocks. A non-stationary random walk has no
278 tendency to return to a fixed value. More generally, mean-reverting series
279 are modelled with first-order autoregression models (ARM) (e.g. Tol and
280 Vellinga [1998], Breusch and Vahid [2008]):

$$281 \quad T_n = a T_{n-1} + S_{n-1} + \epsilon$$

282 Here a is the autocorrelation coefficient, S_n is the deterministic radiative
283 forcing at times n due to any factor: solar variations, increases in greenhouse
284 gases, aerosols, and volcanic eruptions, and ϵ is the random error. We
285 assert that the recurrence equation models not only the errors but also the
286 response of the system itself, particularly when the response time is long.
287 This approach may be too elementary for real system modelling and prone
288 to diagnostic bias Foster et al. [2009] stemming from several decay times,
289 particularly long-term ones (> 100 years) Stern [2005]. A recurrence matrix
290 formula introduces multiple time scales. A physical system with random-
291 walk behavior has no tendency to return to an equilibrium value, effectively
292 having an infinite characteristic decay time. A value of $a < 1$ ensures its
293 eventual return. A good way to assess the behavior of these equations is to
294 examine the response to a constant step forcing S as shown in Fig 10. When
295 the autoregression term a is zero, the response follows the step. If the $a = 1$
296 then the response diverges to plus infinity.

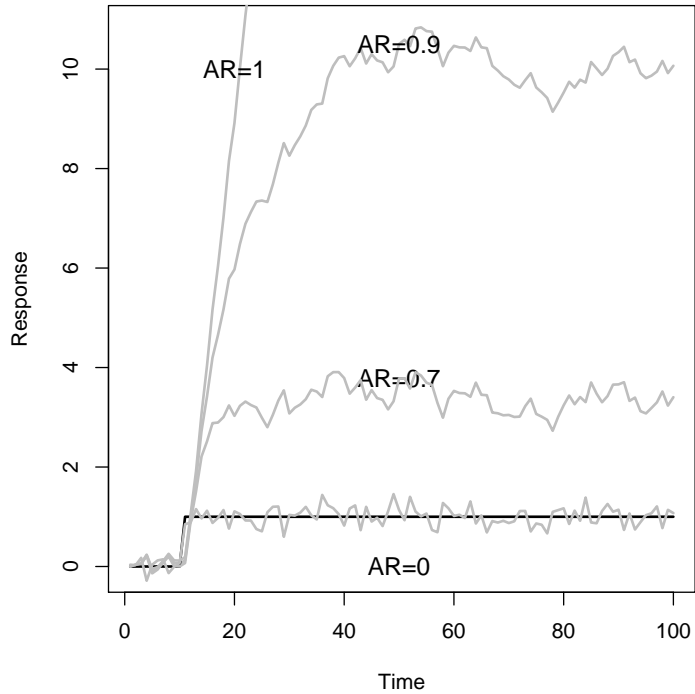


Figure 10: The increasing response to a step forcing (black) of magnitude 1 by recurrence equations with increasing AR value and random noise, for AR values of 0, 0.7, 0.9 and 1.

297 We can calculate the trajectory of a recurrence equation by summing
 298 consecutive terms. The divergent behavior of the random walk in Fig 10 is
 299 due to the ever increasing sum of the constant level forcing S , as follows:

$$300 \quad T_n = T_0 + \sum_{t=0}^n (S + \epsilon_t)$$

301 An equation with AR of less than one (e.g. $a = 0.7$) increases in a
 302 decreasing exponential to a maximum value. A higher AR (e.g. $a = 0.9$)
 303 results in a higher but still limited equilibrium level(Fig 10). The size of this
 304 response to the step forcing is indicative of the intrinsic gain.

305 Simplifying the equation by dropping the ϵ , the pattern of the recurrence
306 equation which after n steps can be expressed as the summation shown
307 below.

$$308 \quad T_n = a(\dots(aT_0 + \frac{S}{C})\dots) + \frac{S}{C} = a^n T_0 + \frac{S}{C} \sum_{t=0}^n a^t$$

309 It is clear that when $0 < a < 1$, as $n \rightarrow \infty$ then $a^n T_0 \rightarrow 0$ and the limit
310 of the geometric series goes to $\tau = \frac{1}{1-a}$. The single parameter τ governs
311 both the rate of the rise and the height of the final equilibrium, so that the
312 characteristic decay time or rise time, the intrinsic gain and the amplification
313 are equivalent. As an example, the 5C range of glacial-interglacial transitions
314 may be simulated by feeding random energy (mean=0, $\sigma = 0.1C$) into a
315 recurrence equation with an autocorrelation coefficient of nearly one.

316 From the analysis above, there is a clear meaning of AR as fractional
317 retention. A system with $a = 1$ represents a ideal accumulator with no
318 losses. The value $1 - a$ is the loss or leakage at each time step. In the case
319 of a real energy balance model, the equation describes an object absorbing
320 heat and increasing temperature in response to radiative forcing, where the
321 loss is proportional to the current temperature.

322 **3.2 Energy Balance Models**

323 We now derive the recurrence model from the solutions to differential equa-
324 tions in a zero-dimensional energy balance model (EBM). A useful concep-
325 tual model of an EBM is a fluid surge tank, used to moderate pressure and
326 flow in hydraulic systems. S is an unregulated input, C a tank capacity

327 with fluid level T , and F a regulated output. When the surge tank receives
328 a sudden inflow, fluid volume in the tank accumulates and the level T rises.
329 Output F increases proportional to T and the fluid level eventually stabi-
330 lizes. The model responds at an impulsive forcing with exponential decay,
331 and rises to a new level T when the forcing persists.

332 The set of equations below describe a standard EBM Spencer and Braswell
333 [2008]: ΔS is the change in short-wave radiation into the Earth's atmo-
334 sphere, ΔF is the change in longwave radiation leaving the atmosphere, and
335 ΔT is the change in global-mean average temperature at the Earth's surface
336 and T :

$$\Delta T = \lambda \Delta F \tag{1}$$

$$C \frac{dT}{dt} = (S - F) \tag{2}$$

$$\frac{dT}{dt} + \frac{T}{\lambda C} = \frac{S}{C} \tag{3}$$

$$T(t) = \frac{e^{-\frac{t}{\tau}}}{C} \int e^{\frac{t}{\tau}} S(t) dt \quad (4)$$

337 The parameter λ in (1) states that temperature and radiative forcing are
 338 proportional, with units of Kelvin per Watts per meter squared ($K/W/m^2$).
 339 The heat capacity C in (2) gives the proportional rate of change of temper-
 340 ature due to radiative forcing, or imbalance (S-F), with units of Joules per
 341 Kelvin (J/K). Substituting equation (1) into (2) gives the first-order ordi-
 342 nary differential equation (3), with a solution (4). The characteristic decay
 343 time, or rise time τ is given by $\tau = \lambda C$.

344 To gain greater familiarity with the solution, Fig. 11 illustrates the be-
 345 haviour of (4) in response to impulse, constant, and fast and slow periodic
 346 forcing.

347 An impulse forcing (Fig. 11A) decays exponentially after the initial surge
 348 due to the negative exponential. Step forcing (Fig. 11B) converges to a sta-
 349 ble, higher value. Pure integration, unaffected by exponential decay, would
 350 ramp up constantly on a step increase (thick dashed line in Fig. 11B). Short
 351 period forcing (Fig. 11C) shows amplification (but less than in Fig. 11B),
 352 and a prominent phase shift of around 90 degrees. Finally, a slower periodic
 353 forcing (Fig. 11D) shows higher amplification due to the longer period of
 354 warming.

355 The model produces amplification and lag, and a maximum proportional
 356 to the size and duration of the forcing and the decay properties τ of the

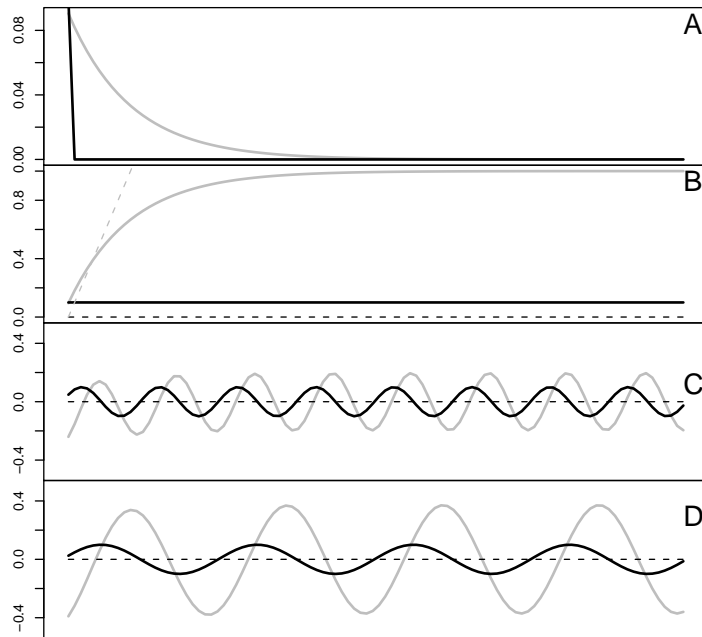


Figure 11: The response (gray) of the energy balance equation 4 to a range of forcings (black) of amplitude 0.1: (A) impulse, (B) step, (C) short periodic and (D) long periodic. Note the apparent amplification of longer period forcing exceeds the short period forcing.

357 system.

358 Equation (3) written as a first order autoregressive model (order=(1,0,0))
359 is:

$$T_t - T_{t-1} + \frac{T_{t-1}}{\tau} = \frac{S_{t-1}}{C} \quad (5)$$

$$T_t = a T_{t-1} + \frac{S_{t-1}}{C} \quad (6)$$

360 Note that $a = 1 - \frac{1}{\tau}$ but the exact discretization is $a = e^{-\frac{\Delta t}{\tau}}$, a source
361 of significant bias only if the rise time is less than the sampling period (e.g.
362 less than a year).

363 3.3 Linked Systems

364 A three-component system with levels atmosphere, surface and ocean basins
365 that accumulate and pass heat between them is as follows:

$$\begin{bmatrix} x_i \\ y_i \\ z_i \end{bmatrix} = \begin{bmatrix} a_x & a_{yx} & 0 \\ a_{xy} & a_y & a_{zy} \\ 0 & a_{yz} & a_z \end{bmatrix} \begin{bmatrix} x_{i-1} \\ y_{i-1} \\ z_{i-1} \end{bmatrix} + \begin{bmatrix} X_{i-1} \\ Y_{i-1} \\ Z_{i-1} \end{bmatrix} \quad (7)$$

366 The diagonal entries in the matrix represent the persistence from one
367 time step to another, while the off-diagonal entries represent the transfer of
368 energy from one level to another. The matrix is in tridiagonal form, where
369 each component communicates with the next, as in a cascade of surge tanks.
370 As the system is at equilibrium, the matrix should also be symmetric which
371 has interpretative and computational benefits.

372 The intrinsic AR coefficients (and decay rates) will match the eigenvalues
373 of the matrix and not the values on the diagonal. In particular, the dominant
374 eigenvalue of the recurrence matrix will be the largest value.

375 4 Parameterization

376 A linear regression can be used to estimate the parameters of the system
377 with the form:

$$378 \quad T_i = aT_{i-1} + bS_{i-1} + c$$

379 where a is the AR coefficient, S is the solar irradiance at the TOA, b is
380 the effect of changes in solar irradiance on global temperature T , and c is
381 the intercept that allows us to calculate the equilibrium value. The result of
382 fitting the HadCRU annual temperature to the solar irradiance from Lean
383 et.al. (2001) are as follows:

$$384 \quad a = 0.89 \pm 0.04; b = 0.063 \pm 0.029C/Wm^2; c = -86.2 \pm 39; R^2 = 0.8603$$

385 The solar effect on temperature is $0.06 \pm 0.03C/Wm^2/yr$. The volume
386 of water that would rise by 0.06C after one year of even heating by 1Watt

387 has a depth of 159 meters, corresponding to the midpoint depth of the
 388 tropical ocean thermocline. We can also calculate the equilibrium value of
 389 $86.2/0.063 = 1365.9Wm^2$ as found previously.

390 4.1 Earth Components

391 A standard ARIMA fitting procedure with order=(1,0,0) estimated the AR
 392 and SD parameters of natural data sets. The AR in the atmosphere (Table 1)
 393 decreases from 0.5 for the lower troposphere, to 0.2 in the upper troposphere.
 394 The AR at the surface (Table 2) is around 0.9. The AR of the EPICA data
 395 is indistinguishable from one due to limitations of the ARIMA algorithm,
 396 while the AR of the Zachos sediment core data Zachos et al. [2001] is lower
 397 than expected, probably due to data gaps. However, an estimated decay
 398 time of 3500 years from the spectral plot and the bandwidth relationship
 399 noted previously would give $a = 0.9997$.

	AR	sd1	SD	tau
UAH	0.4969	0.0276	0.1701	1.99
TLT	0.5427	0.0252	0.1654	2.19
TMT	0.2853	0.0324	0.1584	1.40
TTS	0.1328	0.0464	0.1600	1.15
TLS	0.7685	0.0120	0.2507	4.32

Table 1: Estimates of the AR coefficients and SD of atmospheric temperature data from 1979 to present.

400 We also note in passing that $AR \propto \log(\text{Height})$ indicating a functional
 401 relationship of AR to the density of the atmosphere, and thus to λC and τ .

402 According to the accumulation model, system gain increases with more

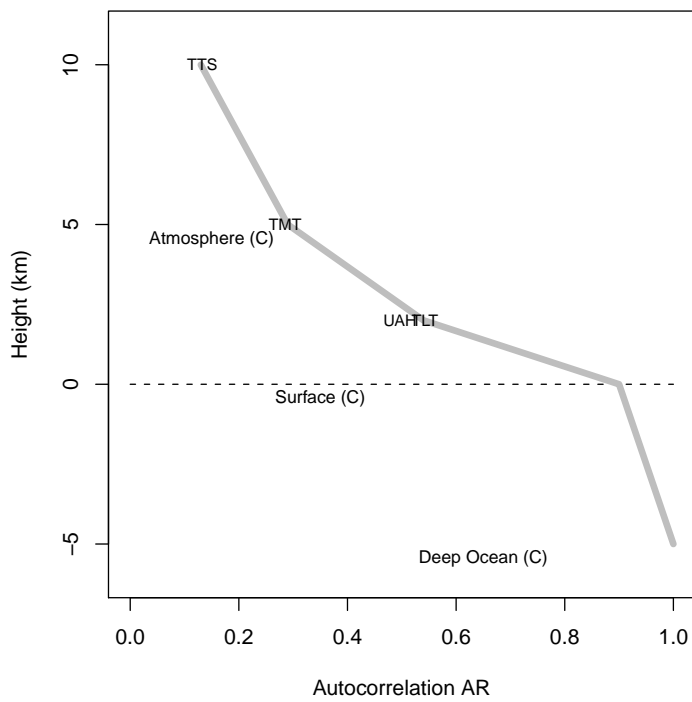


Figure 12: The AR coefficient of the global temperature series decreases with height. This is indicative of a system that loses less energy (retains more energy) per step, with increasing depth.

	AR	sd1	SD	tau
HadCRU	0.9298	0.0009	0.0998	14.25
HadSST	0.9152	0.0010	0.1004	11.80
HadLST	0.8571	0.0016	0.1689	7.00
GISS	0.9068	0.0014	0.1069	10.73

Table 2: Estimates of the AR coefficients and SD of surface temperature data from 1850 to present.

	AR	sd1	SD	tau
vosreg	1.0000	-0.0000	0.7204	480338.72
zachos	0.7415	0.0006	0.2854	3.87

Table 3: Estimates of the AR coefficients and SD of paleoclimate datasets from 800,000BC

403 efficient accumulation of shocks and lower rates of loss. The intrinsic gain
404 increases with decreasing altitude, from 1 or 2 to 10 at the surface and
405 extremely high in the deep ocean (Fig 12). Thus, system response to a
406 forcing depends not only on (1) the size of a forcing, and (2) its duration
407 (affecting the accumulation of heat), but also (3) the forcing depth in a
408 system. For example, long-wave forcing of the low AR, high loss atmospheric
409 level by GHGs would differ from shortwave solar radiation forcing the surface
410 layers of the land and ocean. Geothermal heating in the deep ocean would
411 have the highest intrinsic gain, due to reduced losses.

412 Furthermore, it is clear that the different values for climate sensitivity
413 produced by different studies could easily result by observing of different
414 parts of the system. Studies of short term atmospheric effects, like the
415 Pinatubo eruption, should necessarily provide low sensitivity (e.g. 0.17 to
416 0.20 $C/W/m^2$ Douglass and Knox [2005]), while surface observations would

417 yield higher values Idso [1998]. Higher sensitivity of $2K/(W/m^2)$ would
418 result from observations of surface ocean layers, such as correlative solar
419 estimates Scafetta [2010a], while the highest sensitivity would result from
420 running coupled ocean GCMs over long periods due to the dominating effect
421 of the deep ocean Hansen et al. [2011].

422 4.2 Climate models

423 Evaluation of computer simulations complements analysis of natural tem-
424 perature series Koutsoyiannis et al. [2008], Douglass et al. [2008], Santer
425 et al. [2008]. Table 4) lists the AR and SD for each GCM illustrated graph-
426 ically in Fig 13, along with the natural satellite (red) and surface data sets
427 (green). Most GCMs differ substantially from the natural AR value, but
428 some GCMs are better than others. Models that showed reasonable agree-
429 ment were NCAR1, NCAR2, MIROC3, MRI, and MIUB. Others may be less
430 useful as test-beds of the natural system. These results indicate a subset of
431 models will be more realistic, and a consensus of models will give inferior
432 results.

433 5 Greenhouse gas and other forcings

434 We have not yet excluded a role for CO_2 in the system dynamics, or shown
435 the effects of CO_2 varying independently due to human emissions. That the
436 atmosphere has failed to warm since 2001, over the same ten year period
437 that atmospheric carbon dioxide concentrations have increased by 5% (which

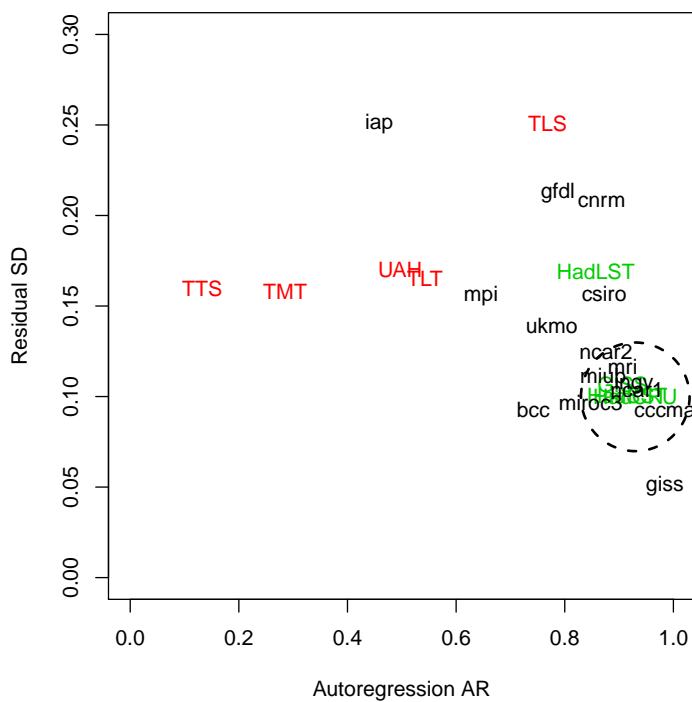


Figure 13: The relative location of GCMs (black), satellite temperature (red) and surface temperature (green) by the AR and SD parameters. The dashed circle encloses the most realistic models.

	AR	sd1	SD	tau
bcc	0.7421	0.0036	0.0930	3.88
cccma	0.9842	0.0002	0.0920	63.20
cnrm	0.8680	0.0019	0.2076	7.58
csiro	0.8724	0.0021	0.1567	7.84
gfdl	0.7873	0.0031	0.2128	4.70
giss	0.9842	0.0002	0.0505	63.12
iap	0.4585	0.0055	0.2509	1.85
ingv	0.9273	0.0012	0.1069	13.76
miroc3	0.8476	0.0020	0.0969	6.56
miub	0.8709	0.0019	0.1116	7.74
mpi	0.6458	0.0042	0.1557	2.82
mri	0.9068	0.0012	0.1165	10.73
ncar1	0.9330	0.0012	0.1039	14.92
ncar2	0.8758	0.0025	0.1247	8.05
ukmo	0.7759	0.0030	0.1393	4.46

Table 4: Estimates of the AR coefficients and SD of general circulation models (GCMs)

438 represents nearly one quarter of all human emissions of carbon dioxide that
439 have occurred since 1751), suggests that the efficacy of the greenhouse gas
440 forcing may be lower than conventionally thought. The accumulative model
441 explains reduced effects from greenhouse gas forcing relative to solar forcing.

442 Atmosphere, surface and deep ocean are forced by longwave, shortwave,
443 and geothermal forcings respectively; x , y and z with forcing S_x , S_y and S_z
444 respectively form a matrix recurrence equation. The eigenvalues of the ma-
445 trix are 0.9997, 0.86 and 0.49, equivalent to intrinsic gains or characteristic

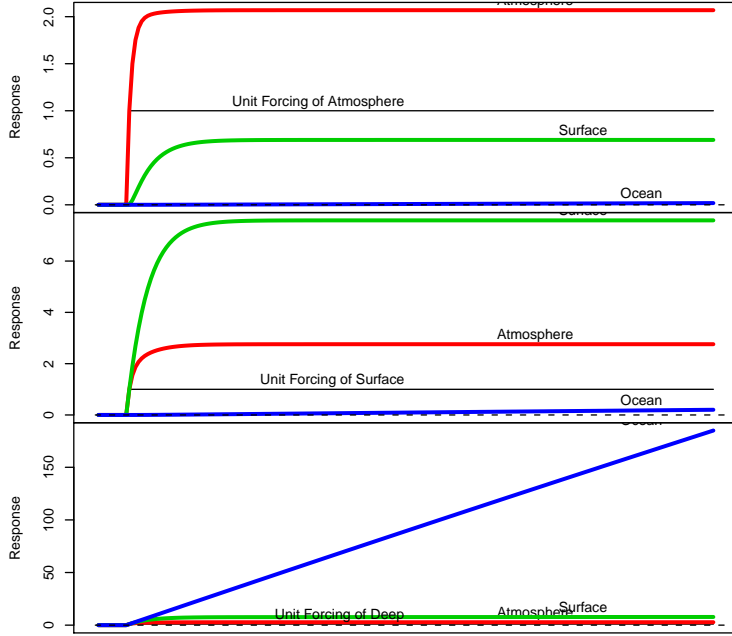


Figure 14: Response to a step forcing (black) in each of the three components: (red) atmosphere, (green) surface and (blue) deep ocean. The step function acts on each component: atmosphere, surface and deep components in the upper, middle and lower panels respectively.

446 rise times of 3300, 7.1 and 2.0 years respectively.

$$\begin{bmatrix} x_i \\ y_i \\ z_i \end{bmatrix} = \begin{bmatrix} 0.5 & 0.05 & 0 \\ 0.05 & 0.9 & 0.00015 \\ 0 & 0.00015 & 0.9997 \end{bmatrix} \begin{bmatrix} x_{i-1} \\ y_{i-1} \\ z_{i-1} \end{bmatrix} + \begin{bmatrix} S_x \\ S_y \\ S_z \end{bmatrix} \quad (8)$$

447 Fig. 14 shows 200 year simulations of the recurrence matrix with step

448 forcing of each level in turn. The range of the responses on the Y axes to an
449 similar step forcing varies by orders of magnitude between atmosphere (x -
450 red), surface (y - green) and deep ocean (z - blue).

451 Forcing of the atmosphere reaches an equilibrium value of two (as $a =$
452 0.5), the surface reaches 0.5 and the deep ocean component is virtually
453 unchanged. Unit forcing of the surface stabilizes at 8 (as $a = 0.86$) while the
454 atmosphere rises to its previous equilibrium value of two. The deep ocean
455 rises slightly. Unit forcing of the deepest component causes the surface and
456 atmospheric components to respond to their equilibrium values and the deep
457 ocean to increase to extremely high levels (as τ is large).

458 The system can be visualized as a downward energy cascade with upward
459 losses. Forcing at the atmospheric, low AR, open end of the cascade, affects
460 deeper components only after losses. In contrast, forcing of deeper compo-
461 nents pushes shallower components to their peak equilibrium response, as
462 the energy moves upward through the cascade. The net result in the exam-
463 ple above is an order of magnitude difference in the effect of a unit forcing
464 at atmosphere and surface of 0.5 and 8 units respectively.

465 The off-diagonal terms in the matrix determine the attenuation, and
466 they are quite uncertain. In particular, the uncertainty of the energy flows
467 across the thermocline is high Wigley [2005] and has recently been estimated
468 at $2 \times 10^{-6} m^2/s$ or $0.75 m^2/year$, 50 times less than the accepted value used
469 in most climate models Douglass et al. [2006a], resulting in a significant
470 overestimate of net human forcing Hansen et al. [2011].

471 The model suggests attenuation of forcing originating in the atmosphere

472 at the surface, and then again at a deeper ocean level. Thus, it is not
473 necessarily the case that warming of the atmosphere by greenhouse gasses
474 would warm the surface or the deep ocean greatly.

475 **6 Discussion**

476 The accumulation theory is supported by a range of evidence. Accumulated
477 solar irradiance has a extremely high ($R^2=0.7$) correlation with global tem-
478 perature since 1950, with an excellent fit to the 11-year solar cycle and trend,
479 but uncorrelated with direct solar radiation (Figs. 4 to 7). The only free
480 variable in the model is the equilibrium value, determined as $1365.9 W/m^2$,
481 resulting in a solar forcing averaging $0.18W/m^2$ at the top-of-atmosphere
482 over the period since 1950. Solar irradiance at TOA is above average for
483 the period Lean [2001], Usoskin et al. [2003], Solanki et al. [2004]. These
484 results are also consistent with previous phenomenological studies attribut-
485 ing more than half of the global temperature change since 1900 Scafetta and
486 West [2007] and 60% of the change since the 1970's Scafetta [2009, 2010a]
487 to natural climate oscillations.

488 Furthermore, the linearity of the spectral frequency plot shows that ac-
489 cumulation is the dominant mechanism of climate change (Fig. 1). There
490 may be a limit to the range of temperature swings at around 22,000 years,
491 suggesting a characteristic decay time of 3500 years from the $2\pi\tau$ bandwidth
492 relationship. It is both surprising and compelling, that annual temperature
493 variance is sufficient to represent global temperature dynamics over these

494 scales, with short characteristic decay times and lags arising from the grad-
495 uated mass density of the atmosphere.

496 The conventional view is that changes in human-caused emissions of
497 greenhouse gases, aerosols and surface albedo cause 20th century warm-
498 ing Duffy et al. [2009]. These same factors support the amplification that
499 cause glacial-interglacial transitions and paleoclimate temperature variation
500 Hansen et al. [2011]. In the alternative theory developed here, changes in
501 temperature due to accumulation of solar heat causes changes in greenhouse
502 gasses and surface albedo.

503 The discussion is structured around the main objections to a large solar
504 influence after Duffy et.al. (2009).

505 The first view (H_0) is that changes in solar radiation have little effect
506 on global temperature, and that changes in greenhouse gas concentrations
507 explain the majority of contemporary and paleoclimatic variability. The
508 second view (H_a) is that solar variation has a significant effect, greater than
509 50% with the residual trend due to other factors such as urban heat island
510 effect, natural cycles, cloud albedo changes or greenhouse gasses. Note that
511 the H_a does not exclude the possibility of observable effects from rising
512 greenhouse gas concentrations, whereas H_0 excludes the possibility of strong
513 solar effects.

514 The first objection to H_a is that the amplitude of the 11 year solar cycle
515 is no more than a few hundredths of a degree. The amplitude would be
516 much larger if solar sensitivity is high North et al. [2004]. We have shown
517 an 11 year cycle with magnitude of more than 0.1C using the accumulated

518 solar variable (Fig. 7). This argument incorrectly assumes that the solar
519 effect must be fast and direct, not slowly integrated.

520 Have variations in TSI have been too small to have contributed to global
521 warming over the last few solar cycles Foukal et al. [2006]? The accumulative
522 model only requires that the Sun's brightness be greater than average over
523 the period, and indeed there has been a long term increase in solar flux that
524 peaked in 1986 with the Grand Solar Maximum Usoskin et al. [2003], Solanki
525 et al. [2004], Lockwood et al. [2009]. We showed that accumulated surplus
526 solar irradiation can explain most of the increase in temperature over the
527 period.

528 Is the hypothesis that solar variability is the dominant climate effect a
529 'non-solution to a non-problem', as direct forcing by albedo, ice extent and
530 vegetation, aerosols, and greenhouse gasses adequately explain temperature
531 change Duffy et al. [2009]? The accumulation theory is more parsimonious,
532 explaining the main features of 20th century warming and the magnitude of
533 variations from one to one million years with a simple, single variable model.

534 Positive feedbacks may help to trap heat Dessler [2010], but the notion of
535 'equilibrium climate sensitivity' is largely unnecessary as the accumulative
536 response is not associated with a variety of materials with specific properties
537 and lags.

538 The H_0 does not exclude an exotic solar influence on climate by regulat-
539 ing more energetic process, such as the influence of gamma ray flux on cloud
540 albedo Svensmark [2007]. Duffy et.al. (2009) find exotic solar effects unlikely
541 because observed global warming requires a forcing of $0.3W/m^2$ at top-of-

542 atmosphere evenly distributed over the ocean and land. We have shown a
543 similar net forcing from accumulated solar anomaly. Slow equilibration is
544 not an exotic mechanism.

545 Has the effect of CO_2 been confirmed by spectral studies of the upper
546 atmosphere? These measurements are subject to large uncertainties and
547 contrary to Harries (2001), recent data show that the expected strong CO_2
548 absorption band in the 700 to 800 cm^{-1} band does not appear in the obser-
549 vations of the difference radiance range between 1970 and 1997 Lu [2010].
550 Moreover, and the H_a does not preclude some noticeable GHG effects of this
551 nature.

552 The H_a contradicts the strong effect of GHGs shown by extensive GCM
553 simulations IPCC [2007]. Climate models underestimate the observed re-
554 sponse to solar forcing Stott et al. [2003], and poorly parametrise ocean
555 mixing parameters overstating the net human-made forcing Hansen et al.
556 [2011]. We have shown both poor parameterization in GCMs, and that sur-
557 face response to atmospheric forcing can be an order of magnitude lower
558 than the same intensity of surface forcing (Fig. 12) due to energy losses
559 down through the system.

560 It could also be said that GCMs already incorporate the slow equilibra-
561 tion effect though coupled ocean-atmosphere simulations. However, many of
562 the parameters in GCMs are virtually unmeasured and involve considerable
563 uncertainty. The range of AR and SD values derived from a representative
564 set of GCM simulations (Fig. 13) is large and generally unrealistic.

565 Another objection from Duffy et.al. (2009) is that dominant forcing by

566 greenhouse gasses already explains the atmosphere of Venus, cooling of the
567 stratosphere, and other phenomena IPCC [2007]. Firstly, the observable fin-
568 gerprints of GHGs may coexist with the H_a , and secondly most atmospheric
569 phenomena are subject to confounding influences e.g. ozone decline has also
570 caused stratospheric cooling.

571 Duffy et.al. (2009) deprecates the 'new science' required to explain phe-
572 nomenological findings of high solar influence. All theories start from a desire
573 to explain phenomena, constrained by fundamental physics. The accumu-
574 lation model is a physically-justified energy balance models with physical
575 parameters including the mass and density of the atmosphere and ocean.
576 The basis of the theory, like all known physics, is the conservation of en-
577 ergy and the appropriate relationship between radiative forcings and the
578 accumulation of heat.

579 The three dimensional atmosphere/surface/ocean recurrence matrix model
580 may be said to be simplistic. Among the advantages of the approach is free-
581 dom to incorporate additional, say, stratospheric, or land/sea components.
582 The model thus satisfies one of the main requirements of a rigorous fore-
583 casting procedure Green and Armstrong [2008] being only as complex as
584 necessary, representing appropriate functional relationships and system de-
585 composition. Its symmetric tridiagonal matrix form is particularly useful
586 for computationally intensive operations such as Monte Carlo sensitivity
587 testing. The study of emergent behaviour of matrix structures would be
588 worthwhile, as would the wider use of control system theory.

589 **7 Conclusions**

590 Contrary to the consensus view, the historic temperature data displays high
591 sensitivity to solar variations when related by slow equilibration dynamics. A
592 variety of results suggest that inappropriate specification of the relationship
593 between forcing and temperature may be responsible for previous studies
594 finding low correlations of solar variation to temperature. The accumulation
595 model is a feasible alternative mechanism for explaining both paleoclimatic
596 temperature variability and contemporary warming without recourse to in-
597 creases in heat-trapping gases produced by human activities. There are no
598 valid grounds to dismiss the potential domination of 20th century warming
599 by solar variations.

600 **8 Appendix 1. Datasets and Software**

601 The sources of data and analysis were as follows:

602 The data were analysed with the R statistical language (<http://www.r-project.org/>) with the supplied ARIMA function. The code is available
603 on request from the author.

605 The EPICA ice core provides was an 800,000 year temperature record
606 and was downloaded from the National Climate Data Center (<ftp://ftp.ncdc.noaa.gov/pub/data/paleo/icecore/antarctica/vostok/deutnat.txt>) then averaged into 1000 year steps.

609 Temperature variations over the last two millennia were from reconstruc-

610 tions using proxy data by Loehle Loehle [2007] ([http://www.ncasi.org/](http://www.ncasi.org/publications/Detail.aspx?id=3025)
611 [publications/Detail.aspx?id=3025](http://www.ncasi.org/publications/Detail.aspx?id=3025)) and by Moberg Moberg et al. [2005]
612 (<http://www.ncdc.noaa.gov/paleo/globalwarming/moberg.html>). These
613 data were aggregated to 20 year steps.

614 The annual global temperature data over the last 150 years (combined
615 HadCRUT , land LST, and sea SST) were downloaded from the Climate Re-
616 search Unit ([http://www.cru.uea.ac.uk/cru/data/temperature/hadcrut3vgl.](http://www.cru.uea.ac.uk/cru/data/temperature/hadcrut3vgl.txt)
617 [txt](http://www.cru.uea.ac.uk/cru/data/temperature/hadcrut3vgl.txt)).

618 The monthly satellite global temperature data for the troposphere tem-
619 perature records (UAH) were downloaded from Remote Sensing Systems
620 and University of Alabama Huntsville ([http://vortex.nsstc.uah.edu/](http://vortex.nsstc.uah.edu/data/msu/t2lt/uahncdc.lt)
621 [data/msu/t2lt/uahncdc.lt](http://vortex.nsstc.uah.edu/data/msu/t2lt/uahncdc.lt)) and troposphere levels (TST, TMT, TLT)
622 and from Remote Sensing Systems ([http://www.ssmi.com/msu/msu_data_](http://www.ssmi.com/msu/msu_data_description.html)
623 [description.html](http://www.ssmi.com/msu/msu_data_description.html)).

624 The solar irradiance data were downloaded from the KNMI Climate
625 Explorer (<http://climexp.knmi.nl>; MONTHLY MEAN TSI: Lean GRL
626 2000). Monthly sun spot numbers were downloaded from the National Aero-
627 nautics and Space Administration ([http://solarscience.msfc.nasa.gov/](http://solarscience.msfc.nasa.gov/greenwch/spot_num.txt)
628 [greenwch/spot_num.txt](http://solarscience.msfc.nasa.gov/greenwch/spot_num.txt)).

629 Annual solar and volcanic forcing, aerosols and well-mixed greenhouse
630 gases were downloaded from the NASA, Goddard Institute of Space Sciences
631 (<http://data.giss.nasa.gov/modelforce/RadF.txt>).

632 The simulations of GCMs were downloaded from KNMI and labeled ac-

633 cording to parent organization as indicated (<http://climexp.knmi.nl>):

634 BCC: itas_bcc_cm1_20c3m_0-360E_-90-90N_n_++a.txt

635 CCCMA: itas_cccma_cgcm3_1_t63_20c3m_0-360E_-90-90N_na.txt

636 CNM: itas_cnrm_cm3_20c3m_0-360E_-90-90N_na.txt

637 CSIRO: itas_csiro_mk3_5_20c3m_0-360E_-90-90N_n_++a.txt

638 GFDL: itas_gfdl_cm2_1_20c3m_0-360E_-90-90N_n_++a.txt

639 GISS: itas_giss_aom_20c3m_0-360E_-90-90N_n_++a.txt

640 NCAR2: itas_ncar2_pcm1_20c3m_0-360E_-90-90N_n_++a.txt

641 UKMO: itas_ukmo_hadcm3_20c3m_0-360E_-90-90N_n_++a.txt

642 NCAR1: itas_ncar1_ccsm3_0_20c3m_0-360E_-90-90N_n_++a.txt

643 MRI: itas_mri_cgcm2_3_2a_20c3m_0-360E_-90-90N_n_++a.txt

644 MPI: itas_mpi_echam5_20c3m_0-360E_-90-90N_n_++a.txt

645 MIUB: itas_miub_echo_g_20c3m_0-360E_-90-90N_n_++a.txt

646 MIROC3: itas_miroc3_2_medres_20c3m_0-360E_-90-90N_n_++a.txt

647 INGV: itas_ingv_echam4_20c3m_0-360E_-90-90N_na.txt

648 IAP: itas_iap_fgoals1_0_g_20c3m_0-360E_-90-90N_n_++a.txt

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